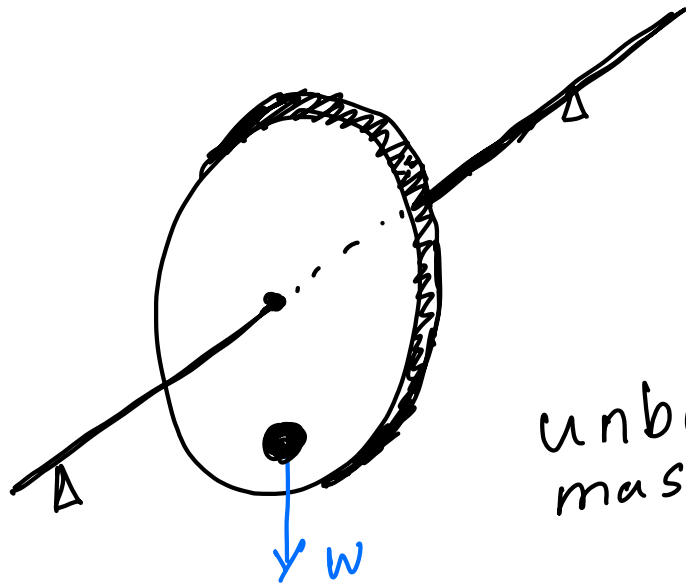


# Rotor Unbalance

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unbalance mass

Static:

Dynamic unbalance:

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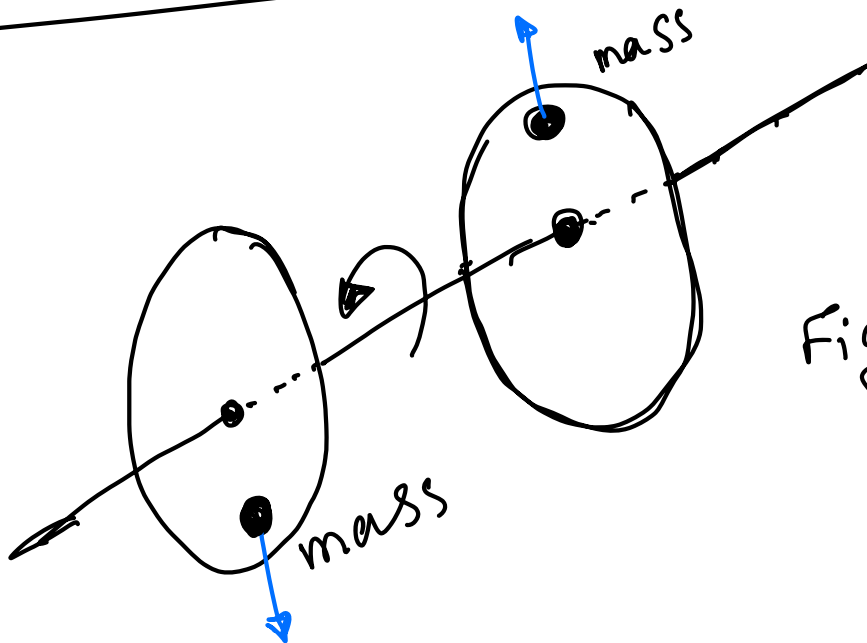
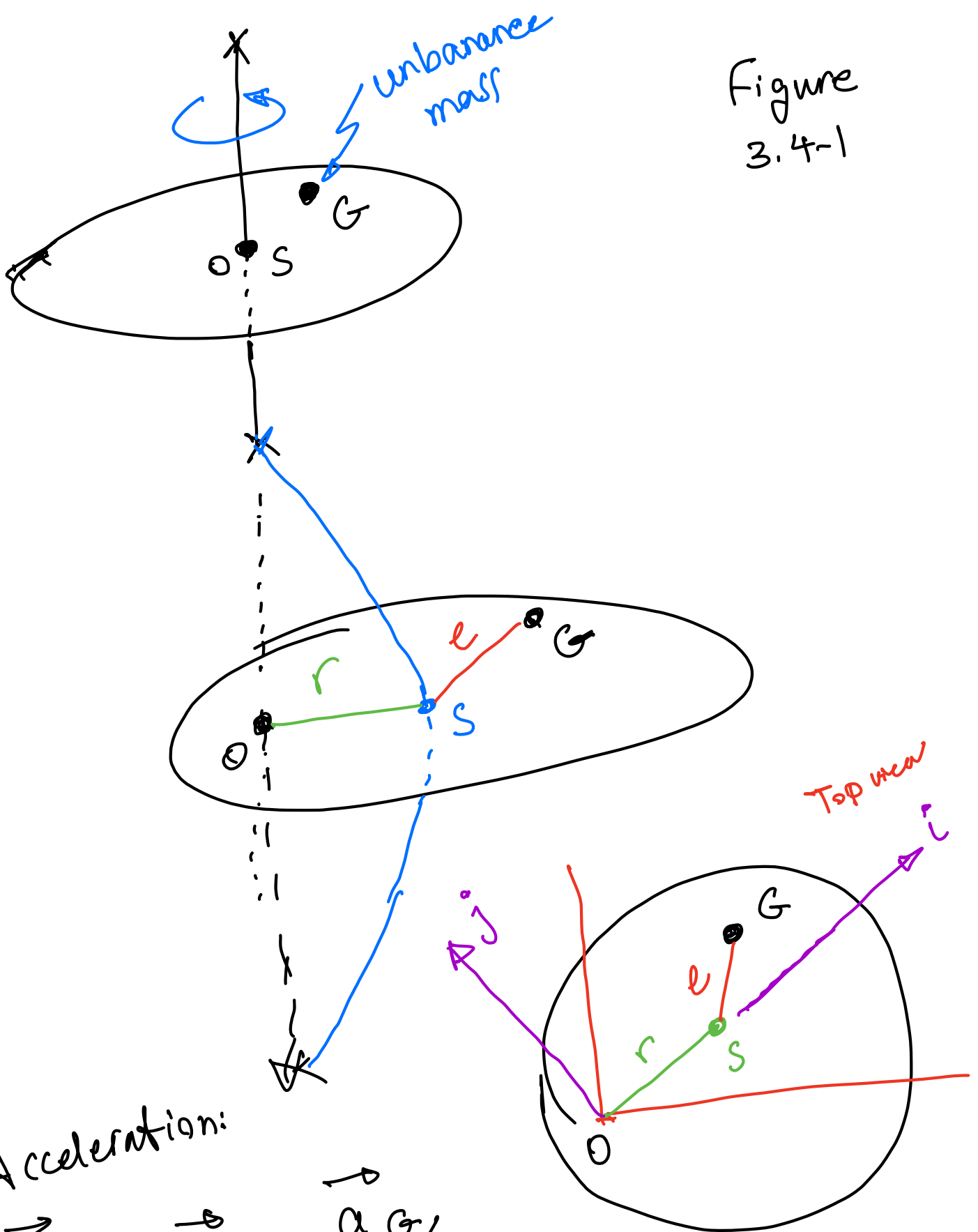


Fig. 3.3-2  
Book

Figure 3.4-1



Acceleration:

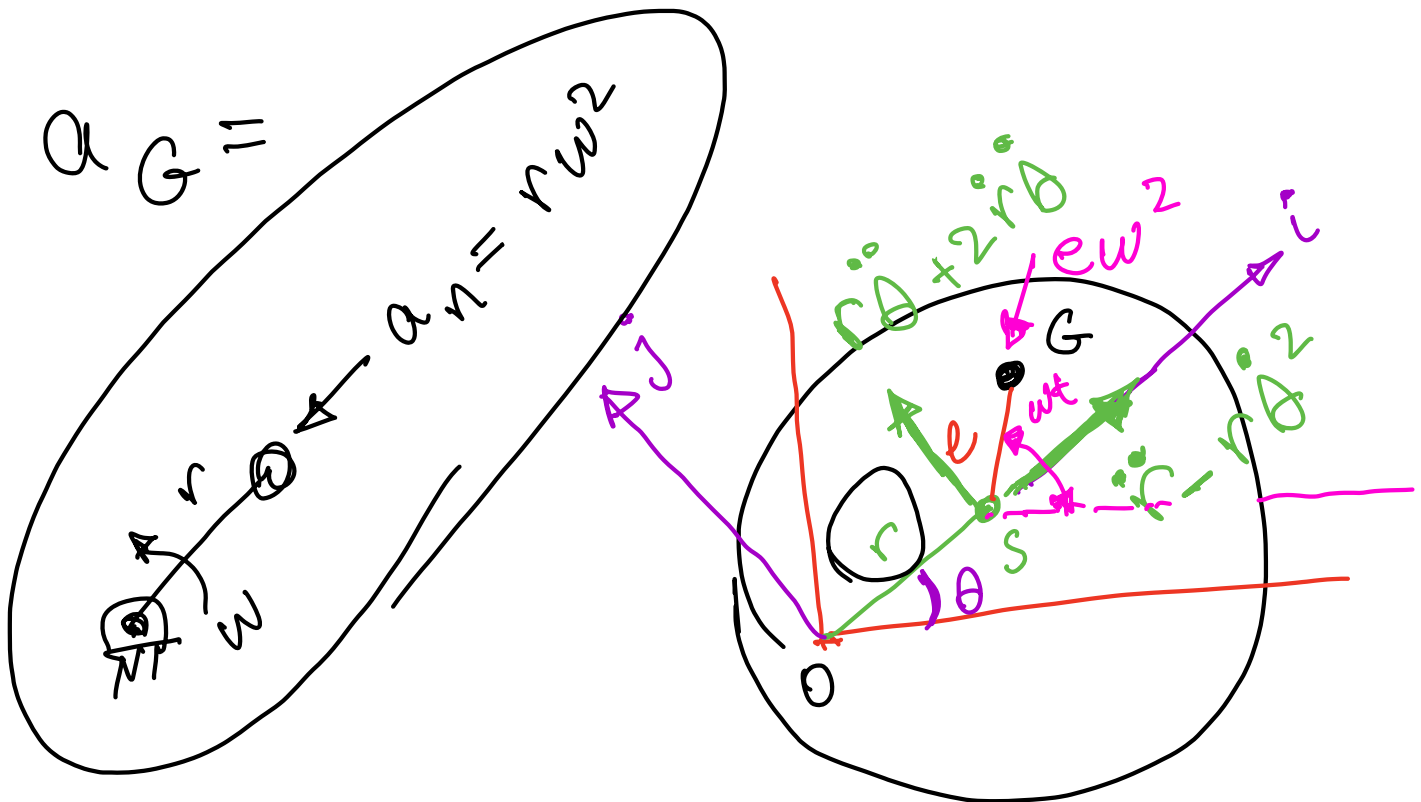
$$\vec{a}_G = \vec{a}_S + a_{G/S}$$

$$\left. \begin{aligned} e &= SG \\ r &= OS \end{aligned} \right\}$$

shaft is rotating at

a constant speed  $\omega$  (the line  $e = so$ )

The line  $r = OS$  whirling at speed  $\dot{\theta}$



$$a_G = [(r\ddot{\theta} - r\dot{\theta}^2) - e\omega^2 \cos(\omega t - \theta)]i$$

$$[(r\ddot{\theta} + 2r\dot{\theta}) - e\omega^2 \sin(\omega t - \theta)]j$$

elasticity of the shaft is with

equivalent stiffness of  $K$   
and equivalent damping of  $c$

$$\boxed{\Sigma F_i = ma_i}$$

$$-kr - c\dot{r} =$$

$$m \left( \ddot{r} - r\dot{\theta}^2 - e\omega^2 \cos(\omega t - \theta) \right)$$

$$\Sigma F_j = ma_j$$

$$-c r \dot{\theta} = m \left( r\ddot{\theta} + 2\dot{r}\dot{\theta} - e\omega^2 \sin(\omega t - \theta) \right)$$

Rearrange:

$$\ddot{x} + 2\zeta \omega_n \dot{x} + \omega_n^2 x = F$$

standard form

$$\ddot{r} + \frac{c}{m} \dot{r} + \left( \frac{k}{m} - \dot{\theta}^2 \right) r = e\omega^2 \cos(\omega t - \theta)$$

$$r\ddot{\theta} + \left( \frac{c}{m} r + 2\dot{r} \right) \dot{\theta} = e\omega^2 \sin(\omega t - \theta)$$

steady state synchronous writ:

$$\dot{\theta} = \omega \Rightarrow \ddot{\theta} = 0$$

$$\ddot{r} = \dot{r} = 0$$

$$\dot{\theta} = \omega \Rightarrow \theta = \omega t - \phi$$

$$\frac{d\theta}{dt} = \omega \rightarrow \int d\theta = \int \omega dt$$
$$\theta = \omega t + \text{constant}$$

$$\left\{ \begin{array}{l} \cancel{r} \ddot{r} + \frac{c}{m} \cancel{r} + \left( \frac{k}{m} - \dot{\theta}^2 \right) r = e\omega^2 \cos(\omega t - \theta) \\ \cancel{r} \ddot{\theta} + \left( \frac{c}{m} r + 2\dot{r} \right) \dot{\theta} = e\omega^2 \sin(\omega t - \theta) \end{array} \right.$$

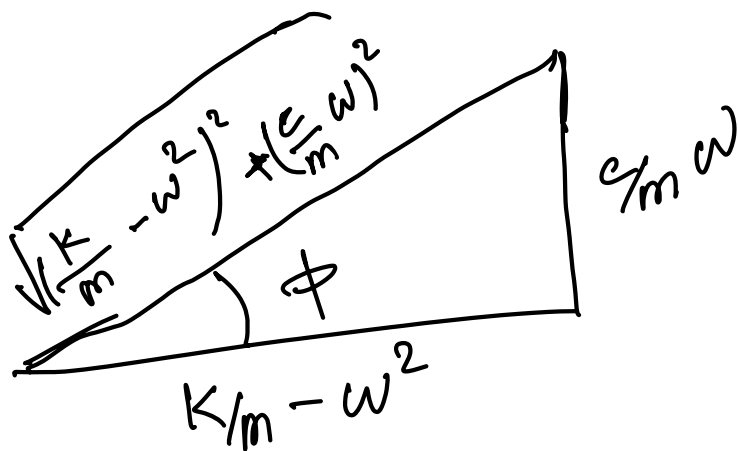
$$\left\{ \begin{array}{l} \left( \frac{k}{m} - \omega^2 \right) r = e\omega^2 \cos \phi \\ \frac{c}{m} \omega r = e\omega^2 \sin \phi \end{array} \right.$$

instead of  $m, c, k$  write the equations in terms of  $\zeta$  and  $\omega_n$ .

The phase angle  $\phi$ :

$$\tan \phi = \frac{\frac{c}{m} \omega}{\frac{k}{m} - \omega^2} = \frac{2 \zeta \frac{\omega}{\omega_n}}{1 - \left(\frac{\omega}{\omega_n}\right)^2}$$

$\omega_n = \sqrt{k/m}$  is the critical speed.  $\zeta = c/c_c$



$$r = \frac{m e \omega^2}{\sqrt{\left(k - m \omega^2\right)^2 + (c \omega)^2}}$$

$$= \frac{e \left( \frac{\omega}{\omega_n} \right)^2}{\sqrt{\left[ 1 - \left( \frac{\omega}{\omega_n} \right)^2 \right]^2 + \left[ 2 \zeta \left( \frac{\omega}{\omega_n} \right) \right]^2}}$$

Example (3.4-1 Book)

Turbines operating above the critical speed must run through dangerous speed at resonance

each time they are started or stopped. Assuming the critical speed  $\omega_n$  to be reached with amplitude  $r_0$

Determine the equation for the amplitude build up with

time. Assume zero damping.

Solution: we will assume synchronous

$$\text{whirl} \rightarrow \dot{\theta} = \omega = \text{constant}$$

$$\ddot{\theta} = 0$$

$$c = 0$$

$$\left\{ \begin{array}{l} \ddot{r} + \left( \frac{k}{m} - \omega^2 \right) r = e\omega^2 \cos \phi \quad (1) \\ 2r\dot{\omega} = e\omega^2 \sin \phi \quad (2) \end{array} \right.$$

$$(2) \Rightarrow \dot{r} = \frac{e}{2} \omega \sin \phi$$

$$\frac{dr}{dt} = \frac{e}{2} \omega \sin \phi \rightarrow \int dr = \int \frac{e}{2} \omega \sin \phi dt$$

$$r = \frac{e\omega}{2} t \sin \phi + r_0 \quad (3)$$

Derivative  $\rightarrow$

$$\dot{r} = \frac{dr}{dt} = 0 \quad (4)$$



①, ③, ④  $\Rightarrow$

$$\left(\frac{k}{m} - \omega^2\right) \left(\frac{e\omega}{2} t \sin\phi + r_0\right) = e\omega^2 \cos\phi$$

constant

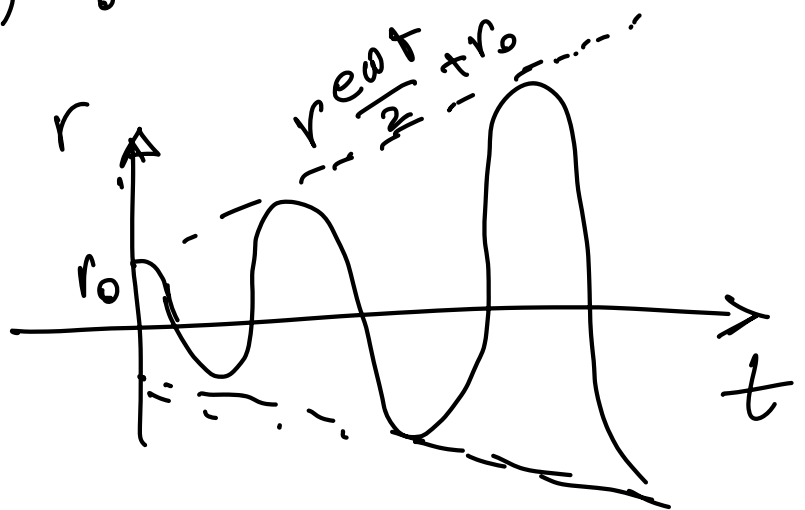
Because the right side is constant then coefficient of  $t$  is zero

$$\left(\frac{k}{m} - \omega^2\right) \sin\phi = 0$$

$\Rightarrow$  Remaining

$$\left(\frac{k}{m} - \omega^2\right) r_0 = e\omega^2 \cos\phi$$

$$\omega = \sqrt{\frac{k}{m}}$$



unstable